#### UNIVERSITY OF WISCONSIN PARTICLE THEORY SEMINAR FEBRUARY 28, 2017



## STERILE NEUTRINO DARK MATTER WITH SUPERSYMMETRY

#### **BIBHUSHAN SHAKYA**





#### BASED ON

#### Sterile Neutrino Dark Matter with Supersymmetry

B. Shakya, J. D. Wells arXiv:1611.01517

#### Cosmological imprints of frozen-in light sterile neutrinos

S. B. Roland, B. Shakya arXiv:1609.06739

#### Sterile neutrino dark matter from freeze-in

B. Shakya

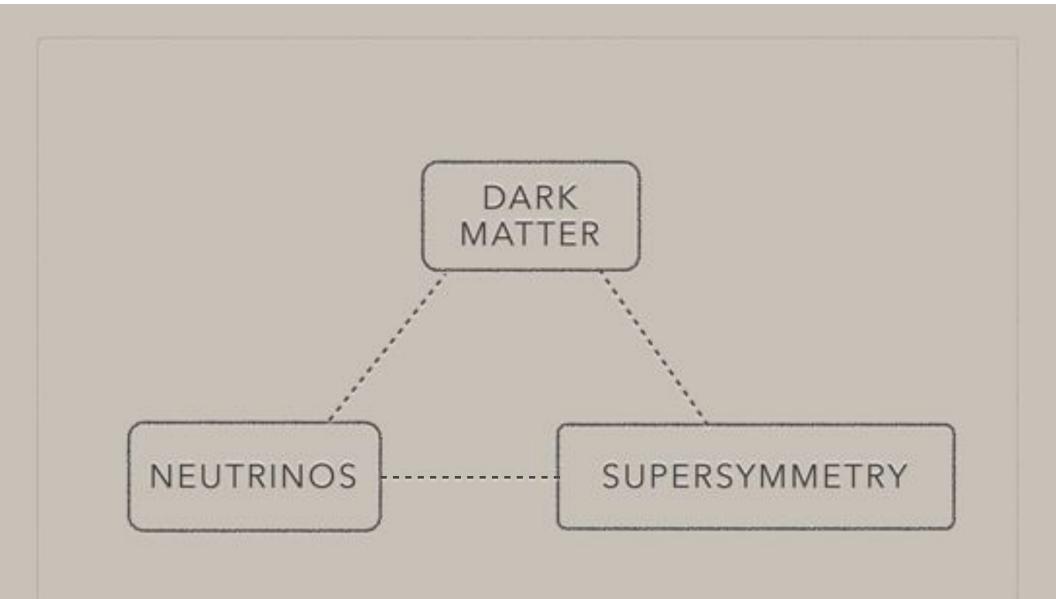
Mod.Phys.Lett. A31 (2016) no.06, 1630005, arXiv:1512.02751

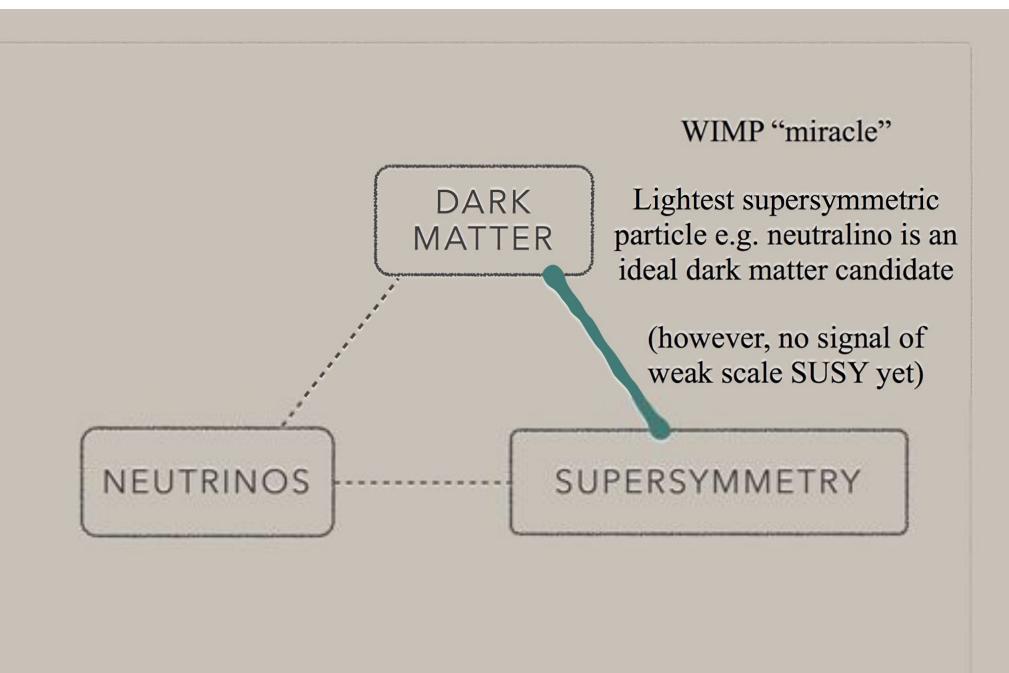
#### PeV neutrinos and a 3.5 keV X-ray line from a PeV scale supersymmetric neutrino sector

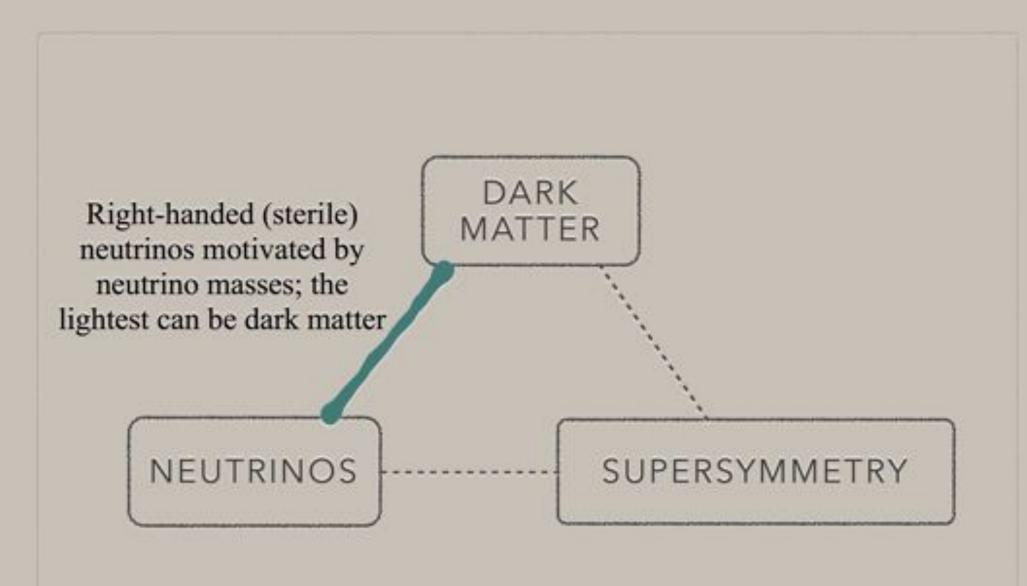
S. B. Roland, B. Shakya, J. D. Wells Phys.Rev. D92 (2015) no.9, 095018, arXiv:1506.08195

#### Neutrino masses and sterile neutrino dark matter from the PeV scale

S. B. Roland, B. Shakya, J. D. Wells Phys.Rev. D92 (2015) no.11, 113009, arXiv:1412.4791







#### 3.5 KEV X-RAY LINE

#### recently in:

#### SEARCHING FOR THE 3.5 KEV LINE IN THE DEEP FIELDS WITH CHANDRA: THE 10 MS OBSERVATIONS

Nico Cappelluti<sup>1,2,3</sup>, Esra Bulbul<sup>4</sup>, Adam Foster<sup>5</sup>, Priyamvada Natarajan<sup>1,2,3</sup>, Megan C. Urry<sup>1,2,3</sup>, Mark W. Bautz<sup>4</sup>, Francesca Civano<sup>5</sup>, Eric Miller<sup>4</sup>, and Randall K. Smith<sup>5</sup> Draft version January 30, 2017

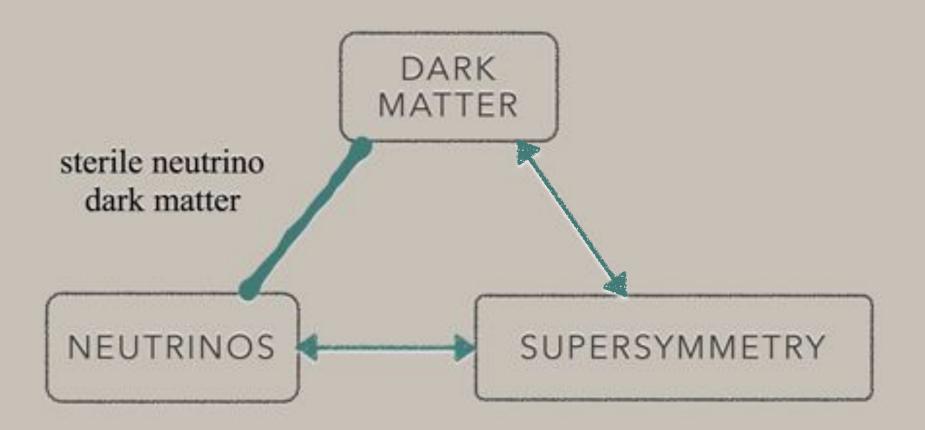
#### ABSTRACT

In this paper we report a  $3\sigma$  detection of an emission line at  $\sim 3.5$  keV in the spectrum of the Cosmic X-ray Background using a total of  $\sim 10$  Ms Chandra observations towards the COSMOS Legacy and CDFS survey fields. The line is detected with an intensity is  $8.8 \pm 2.9 \times 10^{-7}$  ph cm<sup>-2</sup>s<sup>-1</sup>. Based on our knowledge of Chandra, and the reported detection of the line by other instruments, we can rule out an instrumental origin for the line. We cannot though rule out a background fluctuation, in that case, with the current data, we place a  $3\sigma$  upper limit at  $10^{-6}$  ph cm<sup>-2</sup>s<sup>-1</sup>. We discuss the interpretation of this observed line in terms of the iron line background, S XVI charge exchange, as well as arising from sterile neutrino decay. We note that our detection is consistent with previous measurements of this line toward the Galactic center, and can be modeled as the result of sterile neutrino decay from the Milky Way when the dark matter distribution is modeled with an NFW profile. In this event, we estimate a mass  $m_s \sim 7.02$  keV and a mixing angle  $\sin^2(2\theta) = 0.69-2.29 \times 10^{-10}$ . These derived values of the neutrino mass are in agreement with independent measurements toward galaxy clusters, the Galactic center, and M31.

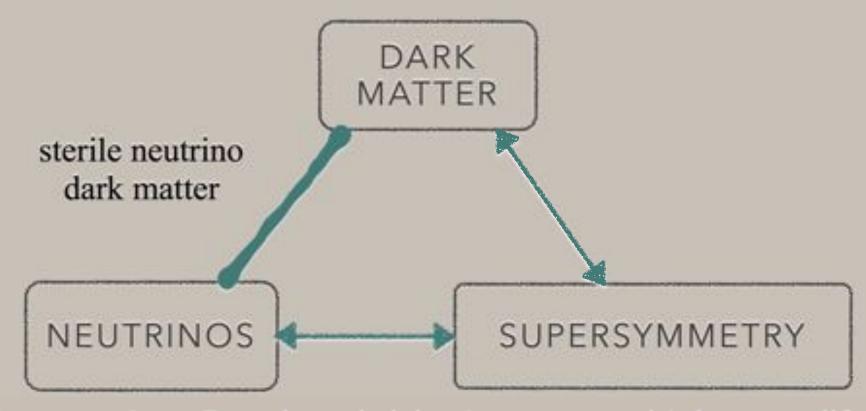
compatible with ~3.5 keV X-ray line from 2014; many papers fitting to ~7 keV sterile neutrino dark matter

potential mismodeling of background (Jeltema and Profumo (2014)), situation unclear; some resolution ~2021?

#### THIS TALK



#### THIS TALK



**Phenomenology:** Does the underlying (supersymmetric) theory modify observable dark matter properties?

**Theoretical:** What can the dark matter properties of the sterile neutrino tell us about the underlying supersymmetric theory?

A QUICK REVIEW

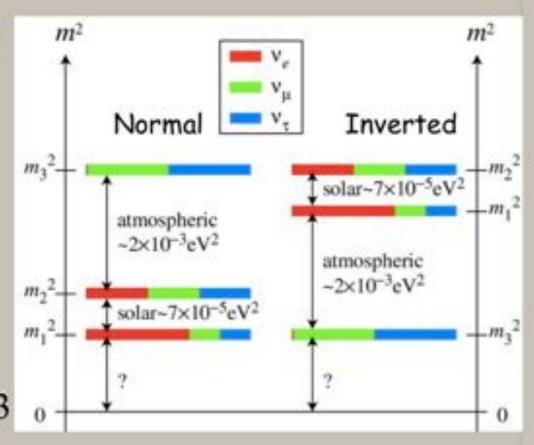
#### **NEUTRINO MASSES**

#### Neutrino sector requires physics beyond the Standard Model!

Neutrinos are massless in the Standard Model

However, solar and atmospheric oscillation data require mass differences.

Absolute mass scale unknown. Measurements constrain the sum to be < 0.23 eV.



#### **NEUTRINO MASSES FROM THE SEESAW**

Add SM-singlet (sterile) right-handed neutrinos Ni

$$y_{ij}L_ihN_j + M_i\bar{N}_i^cN_i$$
 $M \gg y\langle H \rangle \longrightarrow m_a \sim (y\langle H \rangle)^2/M$ 

small mixing between active and sterile states



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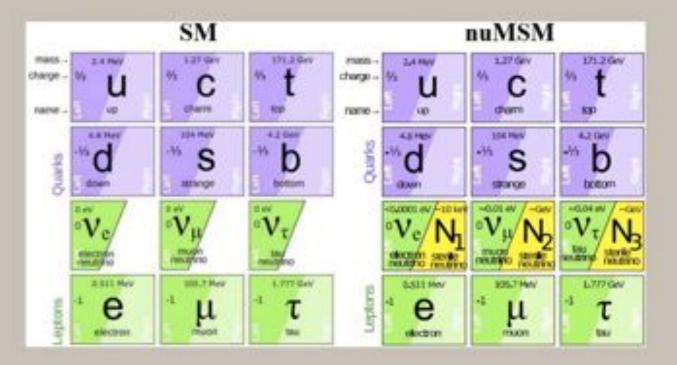
small mixing between active and sterile states

#### What is the mass scale M?

- $y\sim O(1)$ ,  $M\sim 10^{14}$  GeV (GUT scale seesaw)
- $N_1 \sim \text{keV}$  can be dark matter



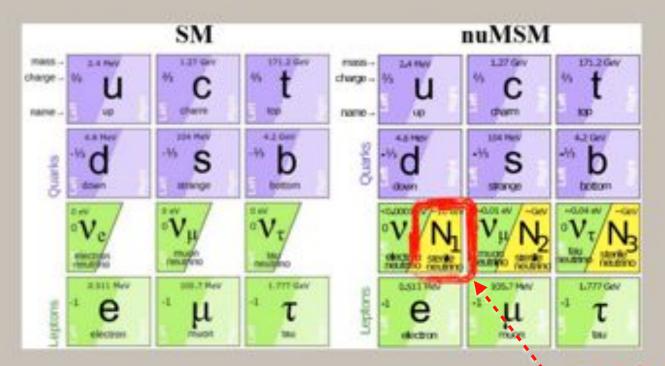
#### Neutrino Minimal Standard Model (vMSM)



extensively studied, explains neutrino masses, baryon asymmetry, and dark matter.

- T. Asaka, S. Blanchet, and M. Shaposhnikov, Phys.Lett. B631, 151 (2005), hep-ph/0503065.
- T. Asaka and M. Shaposhnikov, Phys.Lett. B620, 17 (2005), hep-ph/0505013.
- T. Asaka, M. Shaposhnikov, and A. Kusenko, Phys.Lett. B638, 401 (2006), hep-ph/0602150.

#### Neutrino Minimal Standard Model (vMSM)



#### Dark matter!

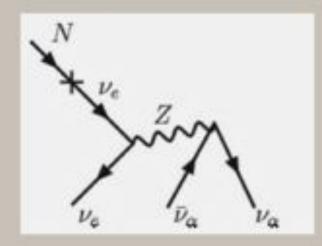
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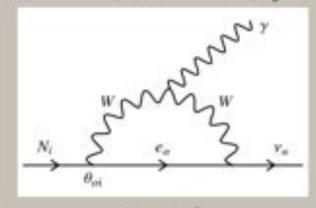
- produced through active-sterile oscillation due to mixing between the two (Dodelson-Widrow mechanism).
- coupling too weak for thermal freeze-out
- · if at keV scale, can be dark matter

$$\Omega_{N_i} \sim 0.2 \left( \frac{\sin^2 \theta}{3 \times 10^{-9}} \right) \left( \frac{m_s}{3 \text{ keV}} \right)^{1.8}$$

lifetime set by 3 body decays

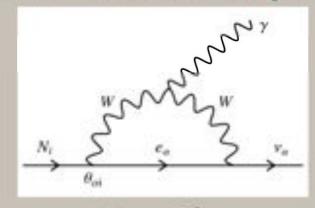


#### Main constraint from X-ray line searches



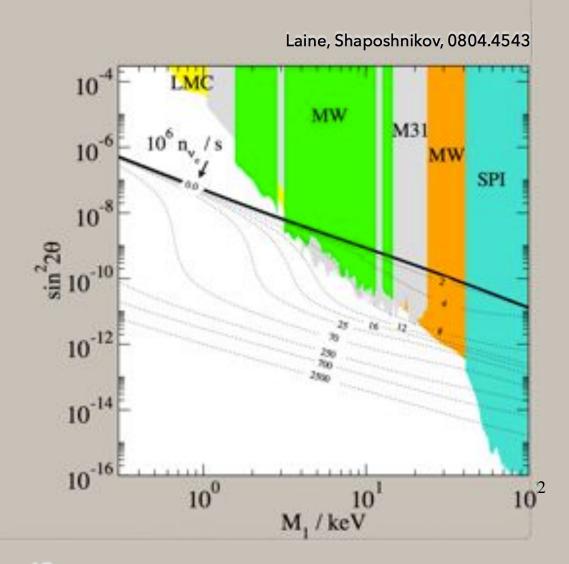
$$\Gamma_{\nu_s \rightarrow \gamma \nu_a} = \frac{9 \alpha_{EM} G_F^2}{256 \cdot 4 \pi^4} \sin^2(2\theta) m_s^5$$

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$$\Gamma_{\nu_s \to \gamma \nu_a} = \frac{9 \, \alpha_{\rm EM} \, G_F^2}{256 \cdot 4 \pi^4} \sin^2(2\theta) m_s^5$$

observations constrain m<sub>v</sub> < O (keV) for sterile neutrino that is all of dark matter



 Lyman-alpha measurements constrain free streaming lengths of warm dark matter candidates (from structure formation)

$$\Lambda_{FS} \approx 1.2 \,\mathrm{Mpc} \left(\frac{\mathrm{keV}}{m_s}\right) \left(\frac{\langle p_s \rangle}{3.15 \, T}\right)_{T \approx 1 \mathrm{keV}}$$

- required to be less than 0.11 Mpc from Lyman-alpha
- SDSS analysis gives m<sub>v</sub> > O(10) keV for production through DW mechanism (eg Viel et al, 0709.0131).

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- SDSS analysis gives m<sub>v</sub> > O(10) keV for production through DW mechanism (eg Viel et al, 0709.0131).

Taken together with the X-ray constraint, rules out sterile neutrino as a dark matter candidate!

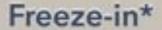
#### Alternatives:

#### Shi-Fuller mechanism

- Presence of lepton chemical potential in plasma can lead to resonantly amplified production of N<sub>1</sub>.
   Colder non thermal distribution, evades Lyman-alpha bounds
- extremely fine-tuned

#### Freeze-out

- additional gauge interactions lead to equilibrium and freeze-out, overproduced abundance fixed by entropy dilution
- potential tension from BBN constraints



\*unrelated to active-sterile mixing

 feeble coupling to some (BSM) particle in the thermal bath leads to gradual production over the history of the Universe

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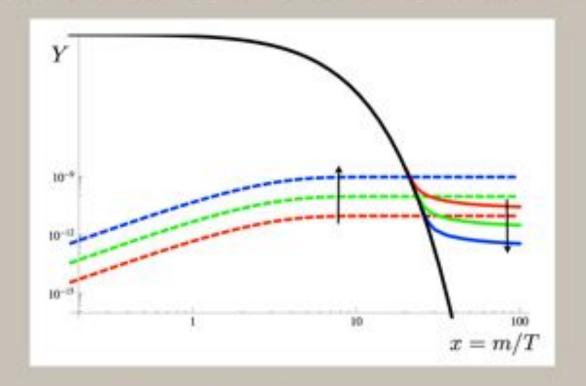
- additional gauge interactions lead to equilibrium and freeze-out, overproduced abundance fixed by entropy dilution
- potential tension from BBN constraints

#### Freeze-in\*

\*unrelated to active-sterile mixing

 feeble coupling to some (BSM) particle in the thermal bath leads to gradual production over the history of the Universe

#### FREEZE-IN VS FREEZE-OUT



from hep-ph 0911.1120

Freeze-out: DM has significant interaction strength with SM, is in equilibrium, decouples

Freeze-in: DM has feeble interaction strength with SM, is never in equilibrium, abundance builds up (freezes in) gradually

many papers; the BSM particle can be [ not an exhaustive list ]

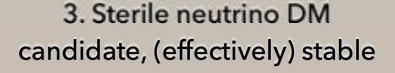
- inflaton (Shaposhnikov, Tkachev, 0604236)
- radion (Kadota, 0711.1570)
- scalar in extended Higgs sector (0711.4646, 0609081, 0702143,1105.1654,1306.3996, 1409.4330, 1411.2773)
- scalar breaking a new symmetry in the neutrino sector (Roland, Shakya, Wells, 1412.4791)

• .......

for a review: Shakya, 1512.02751

#### **Basic ingredients**

 some BSM particle in the early Universe that decays to DM



(technically natural, corresponds to a  $Z_2$  symmetry for  $N_1$ )

[ does not need to be at keV scale ]

2. some feeble coupling ( 
$$x^2 < \frac{m_\phi}{M_{\rm Pl}}$$
 )

$$\mathcal{L} \supset y_{ij} L_i h N_j + x_i \phi \bar{N}_i^c N_i + \lambda (H^{\dagger} H) \phi^2$$

Basic picture



N<sub>1</sub> abundance gradually builds up from phi decays while phi is present in the bath

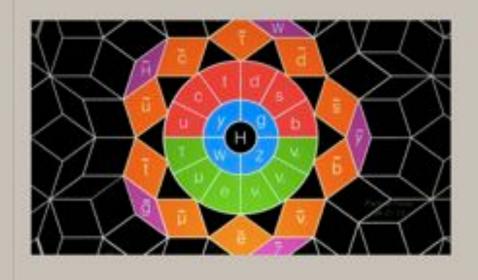
No dependence on active-sterile mixing: can be arbitrarily small, eliminate constraint from X-rays

N<sub>1</sub> produced earlier in the Universe, redshifts and is cooler, more compatible with Lyman alpha

# STERILE NEUTRINO DARK MATTER + SUPERSYMMETRY

## NO HINT OF WEAK SCALE SUPERSYMMETRY.... BUT SURELY IT EXISTS AT SOME SCALE?

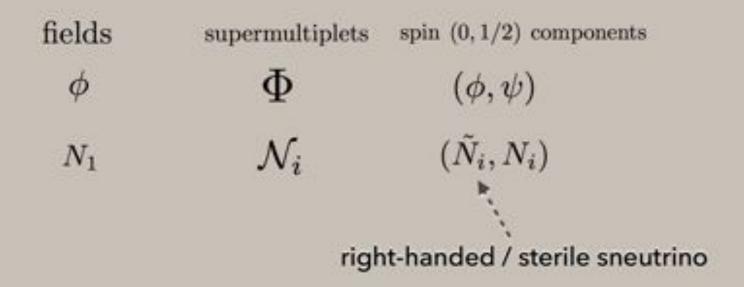
HIERARCHY PROBLEMX



DARK MATTER X

GAUGE COUPLING UNIFICATION

MATHEMATICAL ELEGANCE



<sup>\*</sup> assume R-parity, take stable LSP to be a Higgsino, take to be sub-TeV, forms a small fraction of DM

fields supermultiplets spin (0, 1/2) components

 $\Phi$   $(\phi, \psi)$ 

 $\mathcal{N}_i$   $(\tilde{N}_i, N_i)$ 

Lagrangian  $\mathcal{L} \supset y_{ij} L_i h N_j + x_i \phi \bar{N}_i^c N_i + \lambda (H^\dagger H) \phi^2$ 

comes from superpotential  $W \supset y_{ij} \mathcal{L}_i H_u \mathcal{N}_j + x_i \Phi \mathcal{N}_i \mathcal{N}_i + \sqrt{\lambda} \Phi H_u H_d$ 

in addition to the above Lagrangian (only listing those relevant to this talk):

$$x_i \psi N_i \tilde{N}_i + \sqrt{\lambda} \phi \tilde{H}_u \tilde{H}_d + \sqrt{\lambda} \psi h \tilde{H}$$

$$\mathcal{L}_{soft} \supset y_{ij} A_{y_{ij}} \tilde{L}_i h_u \tilde{N}_j + x_i A_{x_i} \phi \tilde{N}_1 \tilde{N}_1 + \sqrt{\lambda} A_{\lambda} \phi h_u h_d$$

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 $\Phi \qquad \qquad (\phi, \psi)$ 

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possible decay modes for phi

fields supermultiplets spin (0, 1/2) components

 $\phi$   $\Phi$   $(\phi, \psi)$ 

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possible decay modes for sterile sneutrinos (also mixing)

## THE STERILE SNEUTRINO $N_1$

PRODUCTION  $\phi \to \tilde{N}_1 \tilde{N}_1$  if allowed, due to the soft term  $x_i A_{xi} \phi \tilde{N}_1 \tilde{N}_1$  (similarly from psi)

#### **Basic ingredients**

1. some BSM particle in the early Universe that decays to DM 3. Sterile neutrino DM candidate, (effectively) stable

 $\phi$   $N_1$   $N_1$ 

(technically natural, corresponds to  $1 Z_2$  symmetry for  $N_1$ )

 $\mathcal{L} \supset y_{ij} L_i h N_j + x_i \phi \bar{N}_i^c N_i + \lambda (H^{\dagger} H) \phi^2$ 

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- DECAY
- charged under the approximate / exact Z<sub>2</sub> symmetry that stabilizes N<sub>1</sub>.
- Must decay into N<sub>1</sub>; must go through the term  $x_i \psi N_i \tilde{N}_i$ , characterized by the same feeble coupling x<sub>1</sub>. that leads to N<sub>1</sub> freeze-in!

If 
$$m_{\tilde{N}_1} > m_{\psi}$$
,  $\tilde{N}_1 \to \psi N_1$   
if  $m_{\tilde{N}_1} < m_{\psi}$ ,  $\tilde{N}_1 \to N_1 \tilde{H} h$  through an off-shell  $\psi$ 

[ note: both must occur before Higgsino decoupling ]

## THE STERILE SNEUTRINO $N_1$

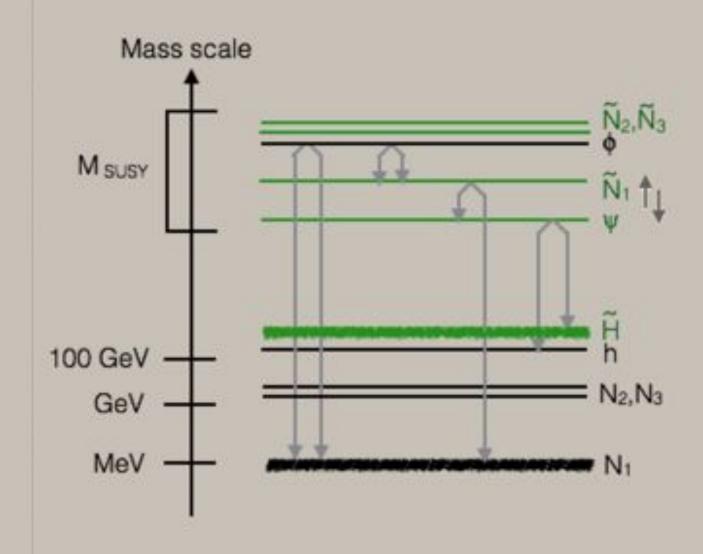
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$$\begin{array}{ll} \text{If } m_{\tilde{N}_1} > m_{\psi}, & \tilde{N}_1 \to \psi N_1 \\ \text{if } m_{\tilde{N}_1} < m_{\psi}, & \tilde{N}_1 \to N_1 \tilde{H} h \end{array} \begin{array}{ll} \text{[ note: both must occur before Higgsino} \\ \text{decoupling ]} \end{array}$$

- each decay produces an N<sub>1</sub> particle
- can be fairly long lived (can even dominate the energy density of the Universe, leading to an early period of matted domination)

### **WORKING SPECTRUM**



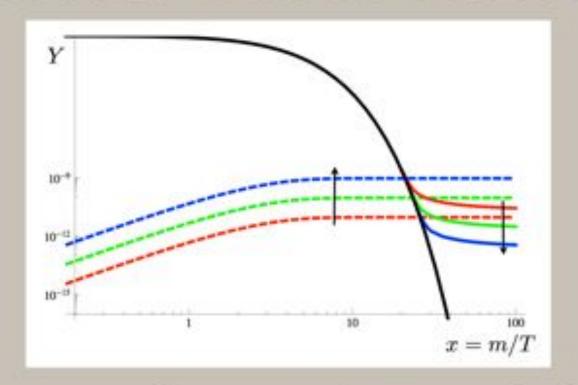
#### RELIC DENSITY AND COMPOSITION

(at least) two distinct production mechanisms

$$\begin{array}{ll} \Omega_{N_1}h^2(\phi) \; \sim \; \frac{10^{24}\,x^2}{2\pi\,S_{N_{2,3}}} \frac{m_{N_1}}{m_\phi} & \text{entropy dilution from} \\ \Omega_{N_1}h^2(\tilde{N}_1) \; \sim \; \frac{10^{24}\,x^2}{2\pi\,S_{N_{2,3}}} \frac{m_{N_1}}{m_\phi} \left(\frac{A_\phi}{m_\phi}\right)^2 & \text{neutrino decays} \end{array}$$

sterile sneutrino decay could provide comparable/ dominant abundance

### FREEZE-OUT VS FREEZE-IN



from hep-ph 0911.1120

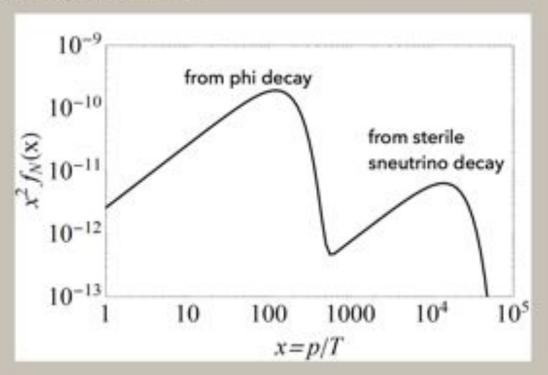
Freeze-out: earlier properties are washed out, decoupling is an IR dominated process

Freeze-in: DM never "thermalizes", final properties are sensitive to details from the early Universe

#### RELIC DENSITY AND COMPOSITION

the two populations don't talk to each other!

second population is hotter (sterile sneutrino is long-lived and decays out of equilibrium)



from hep-ph 1609.06739

extremely nontrivial momentum distribution!

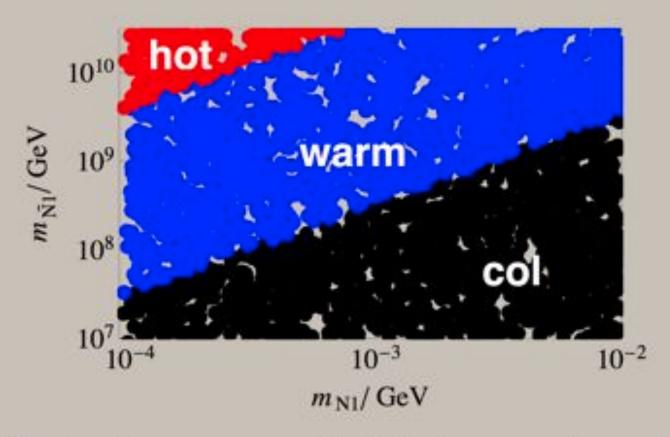
$$\Lambda_{FS} = \int_{t_p}^{t_0} \frac{\langle v(t) \rangle}{a(t)} \, \mathrm{d}t$$

Rough categorization

 $\Lambda_{FS} \lesssim 0.01 \; \mathrm{Mpc}$  cold (most DM models)

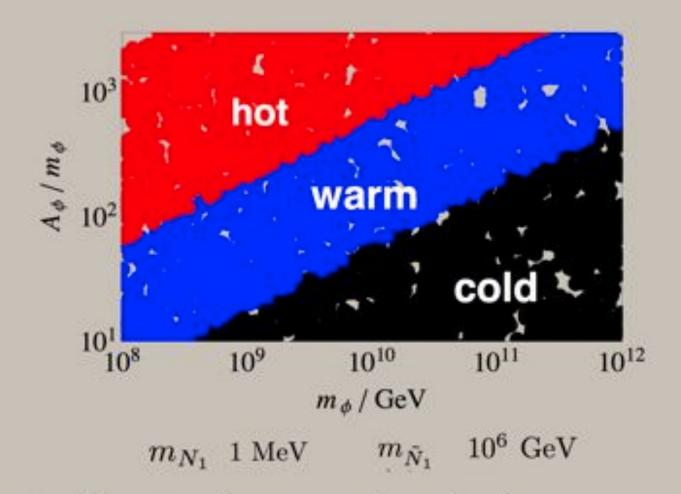
 $0.01 \lesssim \Lambda_{FS} \lesssim 0.1 \; \mathrm{Mpc}$  warm

 $0.1~{
m Mpc} \lesssim \Lambda_{FS}$  hot (ruled out by structure formation)

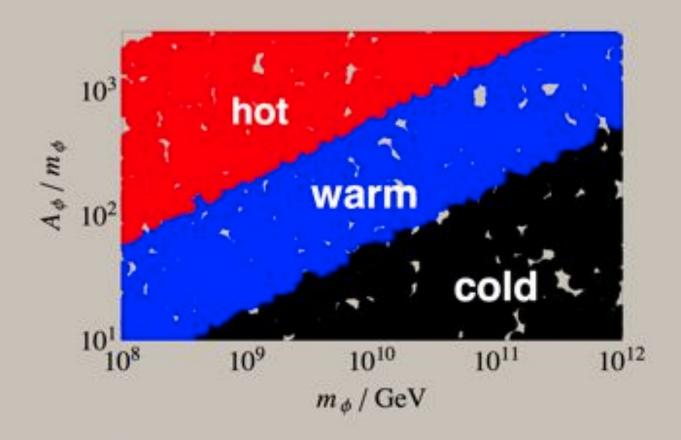


$$ilde{N}_1 
ightarrow \psi N_1 \qquad m_\phi = 10^{11} \; {
m GeV} \qquad A_\phi \, = \, 10 \, m_\phi$$
 coupling x chosen to produce correct relic density

heavier sneutrino / lighter neutrino results in more boosted dark matter



determined not just by masses, but also by parameters that control lifetime!



cold/warm/hot dark matter are all possible in this setup

if the injected population is sufficiently relativistic, can get contributions to the effective number of relativistic degrees of freedom in the early Universe

(at CMB or BBN. in our setup, only BBN)

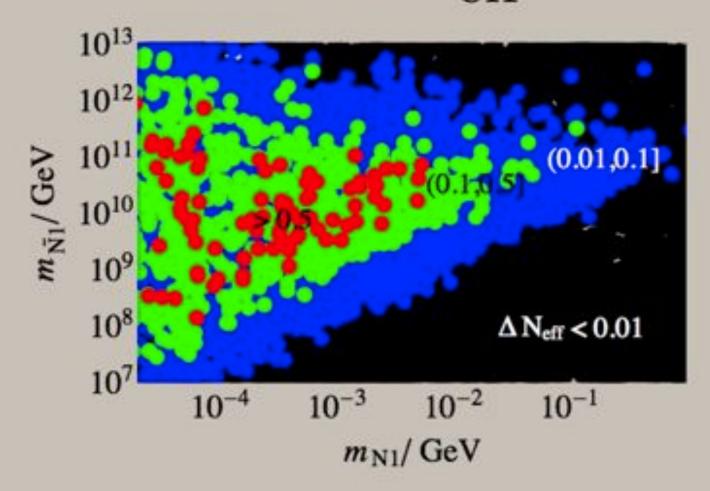
$$\Delta N_{\text{eff}} = \left. \frac{\rho_{N_1}}{\rho_{\nu}} \right|_{T = T_{BBN}}$$

- cannot be all of dark matter, else DM today is too hot and inconsistent with structure formation
- can be a subdominant (<1%) fraction of dark matter, if the rest of dark matter is cold
- generally needs a multi-component dark matter setup; in our framework, N<sub>1</sub> can be both! cold component from phi decay, hot component from sterile sneutrino decay!

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$$\begin{split} \Delta N_{\rm eff} &\approx \frac{10^{-8}}{S_{N_{2,3}}^{1/3} (g_{*SM}/g_{*BBN})^{1/3}} \, \Omega h^2 \frac{m_{\tilde{N}_1}}{T_{decay}} \frac{\rm GeV}{m_{N_1}} \\ &\approx \ 0.2 \left(\frac{\Omega h^2}{0.0012}\right) \left(10^{-8} \frac{m_{\tilde{N}_1}}{T_{decay}}\right) \left(\frac{\rm MeV}{m_{N_1}}\right) \left(\frac{10}{S_{N_{2,3}}}\right)^{1/3} \end{split}$$



 $A_{\phi} = 0.1 m_{\phi}$ 

 $ilde{N}_1 
ightarrow N_1 ilde{H} h$  gives 1% DM. cold population from phi decay makes up the rest

### STERILE NEUTRINO DM

### WITH SUPERSYMMETRY

player in the early Universe; long lived and decays to sterile neutrino DM due to structure of the theory

the sterile sneutrino is an important

- multiple production mechanisms, extends viable parameter space
  - multiple component dark matter with a single constituent
  - can be cold/warm/hot, or some combination of all
  - a subdominant component can give N<sub>eff</sub> contributions, sterile neutrino can still be all of DM

- single production mechanism
- · single component
- can be cold/warm/hot
- cannot be both all of DM and contribute to Neff

#### INVERTING THE QUESTION

SO FAR...

WHAT CAN

SUPERSYMMETRY

TELL US ABOUT

STERILE NEUTRINO DARK MATTER?

#### INVERTING THE QUESTION

NEXT

WHAT CAN

#### STERILE NEUTRINO DARK MATTER

TELL US ABOUT

THE UNDERLYING (SUPERSYMMETRIC) THEORY?

#### Sterile Neutrino Dark Matter from Freeze-In

$$\mathcal{L} \supset y_{ij} L_i h N_j + x_i \phi \bar{N}_i^c N_i + \lambda (H^{\dagger} H) \phi^2$$

#### **ISSUES:**

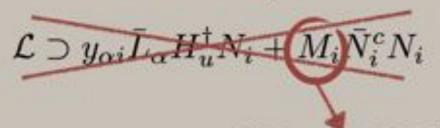
- y~10⁻⁻ to explain neutrino masses
- keV GeV masses for sterile neutrinos put in by hand
- feeble coupling ( $x < 10^{-7}$ ) for DM production

Hints of an underlying structure?

Recall: traditional seesaw requires

$$\mathcal{L}\supset y_{lpha i}ar{L}_{lpha}H_{u}^{\dagger}N_{i}+M_{i}N_{i}^{c}N_{i}$$
 Naively: GUT/Planck scale

Recall: traditional seesaw requires



Naively: GUT/Planck scale

- Assume RH neutrinos charged under a new symmetry: U(1)'
- Prohibits the above terms; traditional seesaw not allowed!

- Introduce an exotic field φ, equal and opposite U(1)' charge to N
- This allows the following terms

$$\frac{y}{M_*}LH_u\mathcal{N}\Phi + \frac{x}{M_*}\mathcal{N}\mathcal{N}\Phi\Phi$$

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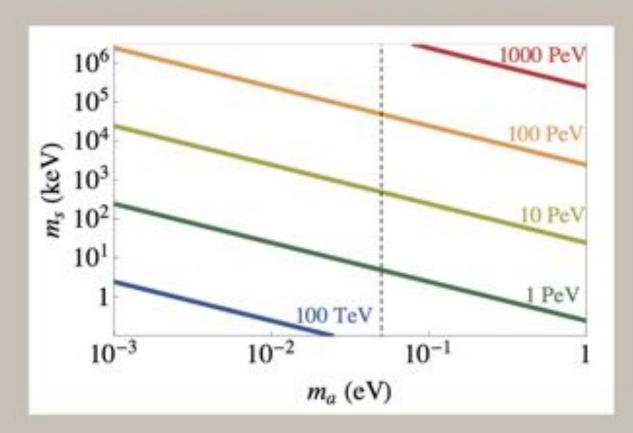
$$\frac{y}{M_*}LH_u\mathcal{N}\Phi + \frac{x}{M_*}\mathcal{N}\mathcal{N}\Phi\Phi$$

If the scalar φ gets a vev, U(1)' broken, effective neutrino mass matrix:

$$M_{\nu} = \begin{pmatrix} 0 & \frac{\langle \phi \rangle \langle H_{u}^{0} \rangle}{M_{\star}} \mathbf{Y} \\ \frac{\langle \phi \rangle \langle H_{u}^{0} \rangle}{M_{\star}} \mathbf{Y}^{\dagger} & \frac{\langle \phi \rangle^{2}}{M_{\star}} \mathbf{X} \end{pmatrix}$$

$$m_s = m_M = \frac{x \langle \phi \rangle^2}{M_*}$$
  $m_a = \frac{m_D^2}{m_M} = \frac{y^2 \langle H_u^0 \rangle^2}{x M_*}$ 

$$\theta \approx \sqrt{\frac{m_a}{m_s}} = \frac{y \langle H_u^0 \rangle}{x \langle \phi \rangle} \qquad m_s = \frac{1}{m_a} \left( \frac{y \langle \phi \rangle \langle H_u^0 \rangle}{M_*} \right)^2$$



Contours of y< $\phi$ >.  $M_* = M_{GUT} (=10^{16} \text{ GeV}), \tan \beta = 2$  0.001 < x < 2

Can get desired active and sterile masses with O(1) couplings and  $<\phi>\sim O(1)-O(100)$  PeV Maps onto vMSM

#### PEV SCALE...SUPERSYMMETRY?

Compatible with m<sub>h</sub>=126 GeV

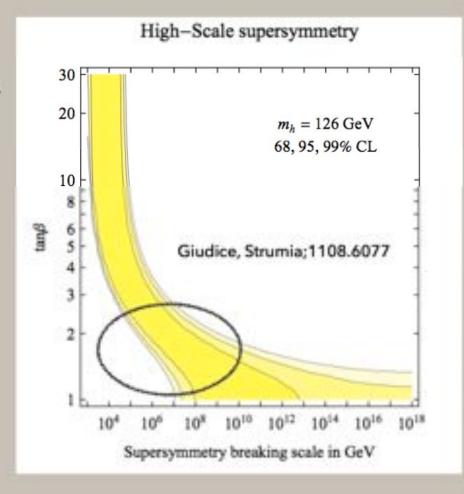
For  $\tan \beta \approx O(1)$ ,  $m_h=126$  GeV implies the scale for supersymmetry (superpartners) is 1-100 PeV

Suggests the vev of φ and the breaking of U(1)' might be related to SUSY breaking.

#### **PeV Scale SUSY**

J. D. Wells (2003), hep-ph/0306127.
 N. Arkani-Hamed and S. Dimopoulos, JHEP 0506, 073 (2005), hep-th/0405159.
 G. Giudice and A. Romanino, Nucl.Phys. B699, 65 (2004), hep-ph/0406088.
 J. D. Wells, Phys.Rev. D71, 015013 (2005), hep-ph/0406088.

ph/0411041.



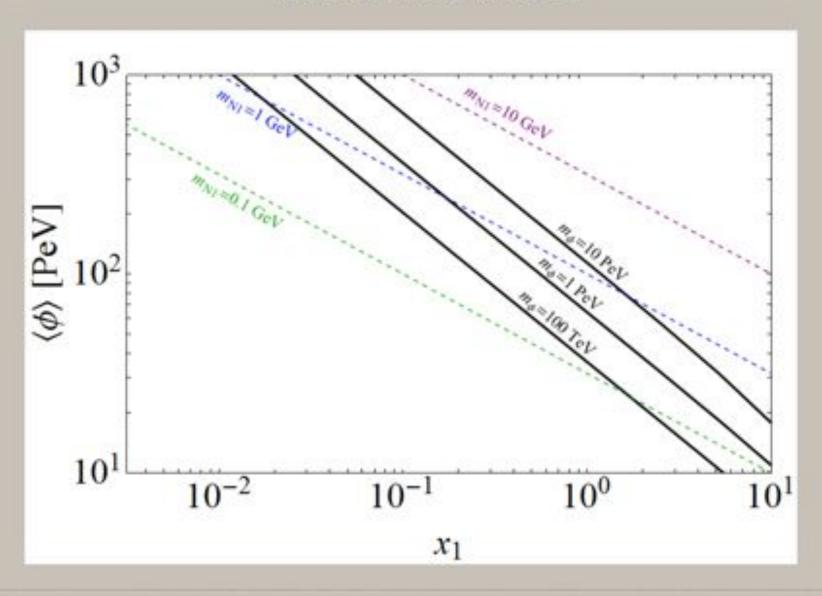
#### STERILE NEUTRINO AS DARK MATTER

#### (Infrared) Freeze-in

$$W \supset \frac{y}{M_*} L H_u \mathcal{N} \Phi + \frac{x}{M_*} \mathcal{N} \mathcal{N} \Phi \Phi$$

Once 
$$\varphi$$
 obtains a vev,  $\phi o N_1\,N_1 \qquad H_u o N_1
u_a \ x_1 = rac{2\,x\,\langle\phi
angle}{M_*} \qquad y_1 = rac{y\,\langle\phi
angle}{M_*} \ \Omega_{N_1}h^2 \sim 0.1\left(rac{x_1}{1.4 imes 10^{-8}}
ight)^3\left(rac{\langle\phi
angle}{m_\phi}
ight)$ 

# PARAMETERS FOR CORRECT ABUNDANCE



#### Sterile Neutrino Dark Matter from Freeze-In

#### **ISSUES:**

y~10<sup>-7</sup> for neutrino masses?

keV - GeV mass scales?

feeble coupling?

**RESOLUTION:**  $\langle \phi \rangle \sim \text{PeV}$ 

$$\sim \frac{\langle \phi \rangle}{M_{GUT}}$$

$$\sim \frac{\langle \phi \rangle^2}{M_{GUT}}$$

freeze-in, small coupling  $\sim \frac{\langle \phi \rangle}{M_{GUT}}$ 

#### STERILE NEUTRINO AS DARK MATTER

#### (Ultraviolet) Freeze-in

$$W \supset \frac{y}{M_*} L H_u \mathcal{N} \Phi + \frac{x}{M_*} \mathcal{N} \mathcal{N} \Phi \Phi$$

(If additional interactions keep  $\varphi$  in equilibrium with thermal bath)

(Doesn't need φ to be in equilibrium)

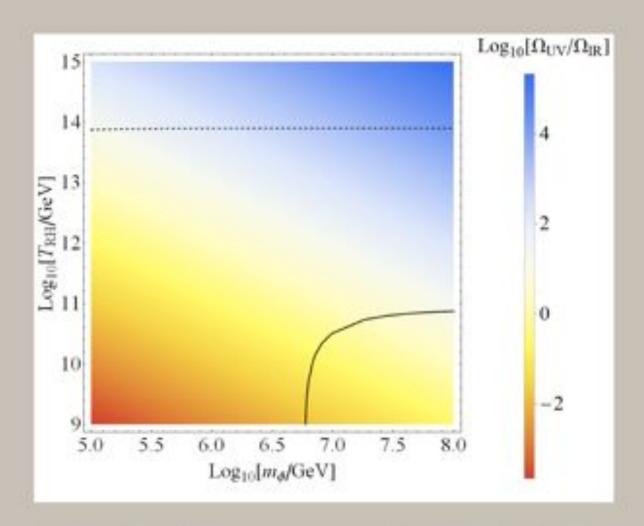
$$\phi \phi \to N_1 N_1$$
,  $\phi H_u \to \nu_a N_1$ ,  $\phi \nu_a \to H_u N_1$ , and  $H_u, \nu_a \to \phi N_1$ 



$$\Omega_{N_1} h^2 \simeq 0.1 \, x^2 \left( \frac{m_s}{10 \, \mathrm{GeV}} \right) \left( \frac{T_{RH} \, M_P}{M_*^2} \right)$$

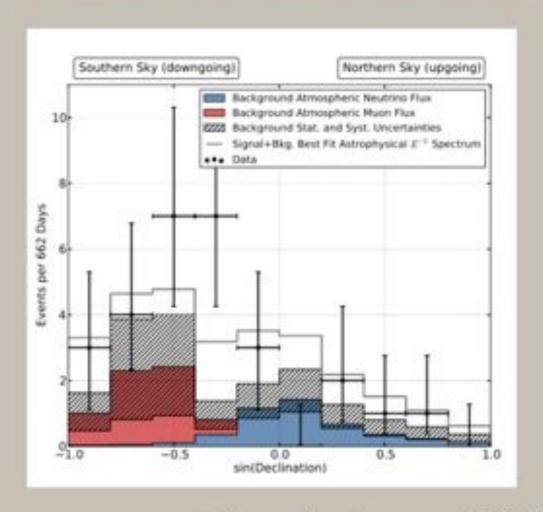
F. Elahi, C. Kolda, and J. Unwin (2014), 1410.6157.
 A. Kusenko, F. Takahashi, and T. T. Yanagida,
 Phys.Lett. B693, 144 (2010), 1006.1731.
 M. Blennow, E. Fernandez-Martinez, and B. Zaldivar,
 JCAP 1401, 003 (2014), 1309.7348.

#### **UV VS IR CONTRIBUTIONS**



dotted/dashed curves : contours of correct relic density

#### **PEV...NEUTRINOS AT ICECUBE!**



37 high energy neutrinos between 30 TeV and 2
 PeV; hint of PeV scale dark matter?

#### **PEV...NEUTRINOS AT ICECUBE!**

Can extend the formalism with the same structure, use the U(1)'
symmetry to stabilize another dark matter component X with a
PeV scale mass that decays to neutrinos [arXiv:1506.08195]

