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## **Noise Voltage and Power Distributions**

## **Noise Voltage**

The voltage V of random noise has a Gaussian probability distribution

$$P(V) = \frac{1}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right),$$

where P(V)dV is the differential probability that the voltage will be within the infinitesimal range V to V + dV and  $\sigma$  is the root mean square (rms) voltage. The probability of measuring *some* voltage must be unity; that is, any probability distribution must be normalized to unity:

$$\int_{-\infty}^{\infty} P(V) dV = 1$$

To confirm the normalization of our noise distribution, we evaluate the integral

$$\int_{-\infty}^{\infty} \frac{1}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right) dV$$
$$= 2 \int_{0}^{\infty} \frac{1}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right) dV$$

$$=\frac{2^{1/2}}{\pi^{1/2}\sigma}\int_0^\infty \exp\left(\frac{-V^2}{2\sigma^2}\right)dV$$

Using the definite integral

$$\int_0^\infty \exp(-a^2 x^2) dx = \frac{\pi^{1/2}}{2a}$$

with  $a^2 = (2\sigma^2)^{-1}$  we get the desired result

$$\int_{-\infty}^{\infty} P(V)dV = \frac{2^{1/2}}{\pi^{1/2}\sigma} \frac{\pi^{1/2} 2^{1/2}\sigma}{2} = 1$$

The rms (root mean square)  $\Sigma$  of a normalized distribution is defined by

$$\Sigma^2 = \langle V^2 \rangle - \langle V \rangle^2$$

For the symmetric Gaussian distribution,  $\langle V \rangle = 0$  so

$$\Sigma^2 = \langle V^2 \rangle = \int_{-\infty}^{\infty} V^2 P(V) dV$$
$$\Sigma^2 = 2 \int_0^{\infty} V^2 \frac{1}{(2\pi)^{1/2} \sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right) dV$$

$$\Sigma^{2} = \frac{2^{1/2}}{\pi^{1/2}\sigma} \int_{0}^{\infty} V^{2} \exp\left(\frac{-V^{2}}{2\sigma^{2}}\right) dV .$$

Using the definite integral

$$\int_0^\infty x^2 \exp(-ax^2) dx = \frac{1}{4a} \left(\frac{\pi}{a}\right)^{1/2}$$

with  $a = (2\sigma^2)^{-1}$  yields

$$\Sigma^2 = \frac{2^{1/2}}{\pi^{1/2}\sigma} \frac{2\sigma^2}{4} 2^{1/2} \sigma \pi^{1/2} = \sigma^2 \ ,$$

demonstrating that  $\sigma$  is really the rms.

## **Noise Power**

A square-law detector multiplies the input voltage V by itself to yield an output voltage  $V_0 = V^2$  that is proportional to the input power. What is the probability distribution  $P_0(V_0)$  of detector output voltage when the input voltage distribution is Gaussian? For simplicity, we set  $\sigma = 1$ . The same value of  $V_0$  is produced by both positive and negative values of V and the probability distribution of V is symmetric, so

$$P_{\rm o}(V_{\rm o})dV_{\rm o} = 2P(V)dV$$

for  $V \ge 0$ . Since  $dV_o = 2VdV$ ,  $dV/dV_o = V_o^{-1/2}/2$  and

$$P_{\rm o}(V_{\rm o}) = \frac{1}{(2\pi)^{1/2}} V_{\rm o}^{-1/2} \exp(-V_{\rm o}/2)$$

for  $0 \le V_o < \infty$ . Notice that the distribution of detector output voltages is sharply peaked near  $V_o = 0$  and has a long exponentially decaying tail, so it looks quite different from a Gaussian distribution.

To confirm that  $P_0$  is properly normalized we evaluate

Planck's sum

$$\int_0^\infty P_{\rm o}(V_{\rm o})dV_{\rm o} = \frac{1}{(2\pi)^{1/2}} \int_0^\infty V_{\rm o}^{-1/2} \exp(-V_{\rm o}/2)dV_{\rm o}$$

using the definite integral

$$\int_0^\infty x^n \exp(-ax) dx = \frac{\Gamma(n+1)}{a^{n+1}} ,$$

where  $\Gamma$  is the Gamma function,  $\Gamma(1/2) = \pi^{1/2}$ , and  $\Gamma(n+1) = n\Gamma(n)$ . Subsituting n = -1/2 and a = 1/2 yields the correct result

$$\int_0^\infty P(V_0) dV_0 = \frac{1}{(2\pi)^{1/2}} \frac{\Gamma(1/2)}{(1/2)^{1/2}} = \frac{1}{(2\pi)^{1/2}} \pi^{1/2} 2^{1/2} = 1 \; .$$

The mean detector output voltage is

$$\begin{split} \langle V_{\rm o} \rangle &= \int_0^\infty V_{\rm o} P_{\rm o}(V_{\rm o}) dV_{\rm o} \\ &= \int_0^\infty \frac{V_{\rm o}}{(2\pi)^{1/2}} V_{\rm o}^{-1/2} \exp(-V_{\rm o}/2) dV_{\rm o} \\ &= \frac{1}{(2\pi)^{1/2}} \int_0^\infty V_{\rm o}^{1/2} \exp(-V_{\rm o}/2) dV_{\rm o} \; . \end{split}$$

Using the definite integral above with n = 1/2 and a = 1/2 yields

$$\langle V_{\rm o} \rangle = \frac{1}{(2\pi)^{1/2}} \frac{\Gamma(3/2)}{(1/2)^{3/2}} \,.$$

 $\Gamma(3/2) = (1/2)\Gamma(1/2) = \pi^{1/2}/2$  so

Planck's sum

$$\langle V_{\rm o} \rangle = \frac{1}{(2\pi)^{1/2}} \frac{\pi^{1/2}}{2} 2^{3/2} = 1 \; .$$

The average detector output voltage is nonzero; it equals the average input power. Had we allowed  $\sigma \neq 1$  we would have gotten  $\langle V_{\rm o} \rangle = \sigma^2$ .

What is the rms  $\Sigma_0$  of the detector output voltage?

$$\Sigma_0^2 = \langle V_o^2 \rangle - \langle V_o \rangle^2$$

so we must evaluate

$$\begin{split} \langle V_{\rm o}^2 \rangle &= \int_0^\infty V_{\rm o}^2 P_{\rm o}(V_{\rm o}) dV_{\rm o} \\ &= \frac{1}{(2\pi)^{1/2}} \int_0^\infty V_{\rm o}^2 V_{\rm o}^{-1/2} \exp(-V_{\rm o}/2) dV_{\rm o} \\ &= \frac{1}{(2\pi)^{1/2}} \int_0^\infty V_{\rm o}^{3/2} \exp(-V_{\rm o}/2) dV_{\rm o} \; . \end{split}$$

Using the definite integral above with n = 3/2 and a = 1/2 yields

$$\langle V_{\rm o}^2 \rangle = \frac{1}{(2\pi)^{1/2}} \frac{\Gamma(5/2)}{(1/2)^{5/2}}$$

 $\Gamma(5/2) = (3/2)\Gamma(3/2) = 3\pi^{1/2}/4$  so

$$\langle V_{\rm o}^2 \rangle = \frac{1}{(2\pi)^{1/2}} 2^{5/2} \frac{3\pi^{1/2}}{4} = 3$$

and

$$\Sigma_{\rm o}^2 = \langle V_{\rm o}^2 \rangle - \langle V_{\rm o} \rangle^2 = 3 - 1 = 2 \; .$$

Thus the rms  $\Sigma_{\rm o} = 2^{1/2}$  of the detector output voltage is  $2^{1/2}$  times the mean output voltage. [If we had kept track of  $\sigma \neq 1$ , we would have gotten  $\Sigma_{\rm o} = 2^{1/2}\sigma^2$ .] The rms uncertainty in each independent sample of the measured noise power is  $2^{1/2}$  times

the mean noise power. If  $N \gg 1$  independent samples are averaged, the fractional rms uncertainty of the averaged power is  $(2/N)^{1/2}$ . This result is the heart of the radiometer equation. According to the central limit theorem, the distribution of these averages approaches a Gaussian as N becomes large.