# A New Dark Matter (in)Direct Search Strategy



**Doojin Kim** University of Wisconsin, WI November 14th, 2017

Based on DK, J.-C. Park, S. Shin, PRL119, 161801 (2017) G. Giudice, DK, J.-C. Park, S. Shin, 1711. xxxxx

# A New Dark Matter (in)Direct Search Strategy at WIMP Detectors



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# Outline

#### I. Introduction/Motivation

Direct detection experiment current status, boosted dark matter search, ...

#### II. Model

Benchmark models, expected signatures, ...

#### **III. Signal Detection**

Benchmark detectors, detection technology, expected signal features, ...

#### **IV. Phenomenology**

Detection prospects, model-independent reach, ...

#### V. Conclusions

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 designed to be
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- $\checkmark\,$  Null observation of WIMP signals
- ✓ A wide range of parameter space already excluded
- ✓ Close to the neutrino "floor"
- ✓ Need new ideas!

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- Overall relic determined by <u>"Assisted" Freeze-out</u> mechanism [Belanger, Park (2011)]
- ✤ Heavier DM  $\chi_0$ : dominant relic, non-relativistic, not directly communicating with SM (hard to detect them due to tiny coupling to SM)
- Lighter DM  $\chi_1$ : directly communicating with SM, subdominant relic (hard to detect them due to small amount)

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 $\Box \chi_1$  can be relativistic at the current universe (non-relativistic as a relic): relativistic DM search

### **Light Boosted DM Detection**

 $\Box$  Flux of boosted  $\chi_1$  near the earth

$$\mathcal{F}_{\chi_1} \sim \frac{\langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1}}{m_0^2}$$
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□ Setting  $\langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1}$  to be ~10<sup>-26</sup> cm<sup>3</sup>s<sup>-1</sup> and assuming NFW DM halo profile, one finds  $\mathcal{F}_{\chi_1} \sim 10^{-7} \text{cm}^{-2} \text{s}^{-1}$  for WIMP mass-range  $\chi_0$ 

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 No sensitivity in conventional dark matter direct detection experiments ⇒ largevolume (neutrino) detectors are motivated, e.g., Super-K/Hyper-K, DUNE



- Elastic scattering [Agashe et al (2014); Berger et al (2014); Kong et al. (2014); Alhazmi et al. (2016)]
- ✓ Inelastic scattering [DK, Park, Shin (2016)]

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Now with GeV/sub-GeV  $m_0 \Rightarrow$  MeV-range  $m_1$  motivated

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Elastic nucleon scattering in the context of gauged baryon number/higgs portal models [Cherry, Frandsen, Shoemaker (2015)]

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- ✓ Expected ER energetic  $\Rightarrow$  MeV sub-GeV range
- ✓ May leave an appreciable track (will be discussed later)
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#### *e*-scattering will be excellent in search for MeV-range (boosted) dark matter particles!



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## **Benchmark Model**

 $-\frac{\epsilon}{2}F_{\mu\nu}X^{\mu\nu}+g_{11}\bar{\chi}_{1}\gamma^{\mu}\chi_{1}X_{\mu}+g_{12}\bar{\chi}_{2}\gamma^{\mu}\chi_{1}X_{\mu}+h.\,c.\,+(others)$  $\mathcal{L}_{int} \exists$ 

- Vector portal (e.g., dark gauge boson scenario) [Holdom (1986)]
- □ Fermionic DM
  - \*  $\chi_2$ : a heavier (unstable) dark-sector state
  - ◆ Flavor-conserving neutral current ⇒ elastic scattering
  - ◆ Flavor-changing neutral current ⇒ inelastic
    scattering [Tucker-Smith, Weiner (2001); Kim, Seo, Shin (2012)]



 $\mathbb{E}$   $\gamma$   $\mathcal{E}$   $\mathcal{E}$   $\mathcal{E}$ 

University of Wisconsin

Hidden



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### **Benchmark Detectors**



Experiment	Geometry	(r, h) or $r$ [cm]	$Mass \ [t]$	Target
XENON1T	Cylinder	(38, 76)	1.0	LXe
DEAP-3600	Sphere	72	2.2	LAr
LZ	Cylinder	(69, 130)	5.6	LXe

[Numbers are for fiducial volumes.]

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#### Dual phase detection technology

Top PMTs
Gas Xe
Liquid Xe
Bottom PMTs

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 Some Xe excited → de-excited, emitting a characteristic scintillation photon (178 nm) detected by PMTs immediately, S1 (scintillation),

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     More Xe ionized, releasing free electrons moving upward by the Drift Field and hitting gaseous Xe,

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     Gaseous Xe excited → de-excited, emitting a

photon detected by PMTs, **S2** (ionization).

#### Dual phase detection technology



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detected by PMTs immediately, **S1** (scintillation),

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  - upward by the **Drift Field** and hitting gaseous Xe,
- 3) Gaseous Xe excited  $\rightarrow$  de-excited, emitting a

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□ Time difference between S1 and S2 giving the depth of the scattering point (~0.1mm resolution)

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### **Detection Technology: XY Plane**

#### □ LOW energy source (<sup>241</sup>AmBe)


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### □ LOW energy source (<sup>241</sup>AmBe)



□ Likelihood analysis allowing position resolution in XY plane as good as < 2 cm (may be better

with high energy source [LUX collaboration (2017)])

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□No dedicated detector studies with highenergetic recoil signals

Doing our best to make as reasonable estimate and expectation as possible

# **High-energetic DM Signal Detection**

□ Point-like scattering position?

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## **High-energetic DM Signal Detection**

□ Point-like scattering position? → Expect a **sizable track**!



[Material property available at NIST (https://physics.nist.gov/PhysRefData/Star/Text/ESTAR.html)]

## **High-energetic DM Signal Detection**

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[Material property available at NIST (https://physics.nist.gov/PhysRefData/Star/Text/ESTAR.html)]

□ Expect tracks of 2 – 10 cm (with LXe) for energy regime of interest

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## **Expected Pattern: Vertical Track**

### A given vertical track

	 I	
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Drift field	ł	

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- 1) can be considered as an array of scattering points,
- 2) Free electrons released at each point: more (less) electrons at the starting (ending) point,
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- 3) Expect a series of flickerings of a few PMTs by an interval of ~10 ns (1 cycle of charge discharge)
- Expect (relatively) easy identification of a lengthy track plus more precise track/energy reconstruction (than the horizontal track in the next slide)

### □ For a given horizontal track

Drift fie	eld		

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 Expect (almost) simultaneous charging of several PMTs, some of which may saturate



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Expect identification of a lengthy track is doable/ achievable Track/energy recon. may require likelihood analysis with unsaturated PMTs

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### □ For a given horizontal track



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# **Positron Signature: Bragg Peak**

### □ A given positron track



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# **Positron Signature: Bragg Peak**

### A given positron track



- 1) stops and gets annihilated with a (nearby) electron,
   creating a characteristic signature of Bragg
   Peak!!!
  - ⇒ Additional handle to identify positrons (or positron tracks)
  - ⇒ Cf.) DEAP having better acceptance for the
     Bragg peak due to its spherical geometry

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# **Expected DM Signals: XY Plane-view**

□ Tracks **POP UP** inside the fiducial volume, **NOT** from outside!



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# **Expected DM Signals: XY Plane-view**



□ Multiple tracks/displaced vertex necessary only for post-discovery (e.g., elastic vs. inelastic)

Cf.) DEAP3600: displaced vertex  $\geq$  6.5 cm identifiable with S1 only by likelihood methods

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□ Any SM backgrounds creating an electron recoil track appearing inside the fiducial volume?

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 $\Rightarrow$  Yes, solar neutrinos, in particular, induced by <sup>8</sup>B.



TABLE II. <sup>8</sup>B neutrino scattering cross sections. The scattering cross sections for <sup>8</sup>B solar neutrinos incident on electrons are given for different values of the minimum accepted kinetic energy  $T_{\min}$ . The neutrinos are assumed to be pure electron neutrinos ( $v_e$ ) or muon neutrinos ( $v_\mu$ ) when they reach the Earth. The cross sections were calculated for  $\sin^2\theta_W = 0.23$ . The quantities  $F_{e,v_\mu}$  and  $F_{e,v_\mu}$  are the fractional changes in the cross section for a change in  $\sin^2\theta_W$  equal to 0.01 [see Eq. (22)].

$T_{\min}$ (MeV)	$(10^{-46} \text{ cm}^2)$	F <sub>e-ve</sub>	$\sigma_{e-\nu_{\mu}}$ (10 <sup>-46</sup> cm <sup>2</sup> )	$F_{e-v_{\mu}}$
0.0	$6.08 \times 10^{2}$	0.029	1.04×10 <sup>2</sup>	-0.040
1.0	$5.09 \times 10^{2}$	0.029	$8.39 \times 10^{1}$	-0.046
2.0	$4.15 \times 10^{2}$	0.028	6.63×10 <sup>1</sup>	-0.052
3.0	$3.27 \times 10^{2}$	0.028	$5.10 \times 10^{1}$	-0.056
4.0	$2.48 \times 10^{2}$	0.028	$3.79 \times 10^{1}$	-0.060
5.0	$1.80 \times 10^{2}$	0.028	$2.71 \times 10^{1}$	-0.063
6.0	$1.23 \times 10^{2}$	0.027	$1.83 \times 10^{1}$	-0.065
7.0	$7.90 \times 10^{1}$	0.027	$1.16 \times 10^{1}$	-0.067
8.0	$4.64 \times 10^{1}$	0.027	$6.76 \times 10^{0}$	-0.068
9.0	$2.44 \times 10^{1}$	0.027	$3.53 \times 10^{0}$	-0.069
10.0	$1.10 \times 10^{1}$	0.027	$1.58 \times 10^{0}$	-0.070
11.0	$3.93 \times 10^{0}$	0.027	$5.64 \times 10^{-1}$	-0.070
12.0	$9.88 \times 10^{-1}$	0.027	$1.41 \times 10^{-1}$	-0.071
13.0	$1.36 \times 10^{-1}$	0.027	$1.94 \times 10^{-2}$	-0.071
13.5	$3.60 \times 10^{-2}$	0.027	$5.13 \times 10^{-3}$	-0.071
14.0	$7.4 \times 10^{-3}$	0.027	$1.0 \times 10^{-3}$	-0.071

[Rev. Mod. Phys., Vol. 59, No. 2, April 1987]

## **Potential Backgrounds**

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□ Estimate only ~0.1 events even at LZ-5yr with an energy cut of  $\geq$  10 MeV (Energy resolution

at  $E_{\text{recoil}} = 10 \text{ MeV}$  is expected to be  $\mathcal{O}(10\%)$  [private communications with experimentalists].)

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### **Benchmark Studies**



FIG. 2: Expected energy spectra of the primary (upper-left panel) and secondary (upper-right panel)  $e^-$  and/or  $e^+$  for four reference points whose details are tabulated in the lower panel.  $g_{12}$  is set to be unity and all mass quantities are in MeV.

 $\chi_2 \text{ long-lived}$   $\ell_{2,\text{lab}} = \frac{c\gamma_2}{\Gamma_2} \sim 16.2 \text{ cm} \times \left(\frac{10^{-3}}{\epsilon}\right)^2 \times \left(\frac{1}{g_{12}}\right)^2$   $\times \left(\frac{m_X}{30 \text{ MeV}}\right)^4 \times \left(\frac{10 \text{ MeV}}{\delta m}\right)^5 \times \frac{\gamma_2}{10}$ 

Fixed Two-body decay of  $\chi_2$  (no displaced vertex)

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Quite <mark>energetic</mark> ER and secondary signals as expected

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### **Benchmark Studies: Detection Prospects**

0.	eliminary		ref	L	rei	f2	ref	3	$\mathrm{ref}$	4
<b>K</b> 1	Expecte	d flux	610	)	43	3	0.9	)8	0.2	4
	Experiments	Run time	$\operatorname{multi}$	single	$\operatorname{multi}$	single	$\operatorname{multi}$	single	$\operatorname{multi}$	single
	VENON1T	1yr	2000	160	220	7.5	0.37	0.37	0.27	0.27
	ALMONTI	$5 \mathrm{yr}$	390	32	43	1.5	0.075	0.075	0.054	0.054
	DEAP-3600	$1 { m yr}$	450	63	55	3.1	—	0.16	_	0.11
		$5 \mathrm{yr}$	91	13	11	0.61	—	0.031	—	0.022
	LZ	$1 { m yr}$	180	27	25	1.3	0.067	0.067	0.048	0.048
		$5 \mathrm{yr}$	36	5.4	5.0	0.26	0.013	0.013	0.0096	0.0096

TABLE II: Required fluxes of  $\chi_1$  in unit of  $10^{-3}$  cm<sup>-2</sup>s<sup>-1</sup> with which our reference points get sensitive to the benchmark experiments. For comparison expected fluxes are shown under the assumptions of  $\langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1} = 5 \times 10^{-26}$  cm<sup>3</sup>s<sup>-1</sup> and the NFW DM halo profile.

□ Selection criteria: "multi" channel – multiple tracks, "single" channel - > 1 track or a single

track with  $E_{\text{recoil}} \ge 10$  MeV.

□ 3 signal events under the zero background assumption.

DEAP3600 having no sensitivity to ref3 and ref4 in the "multi" channel: no displaced vertices

in ref3 and ref4, it is challenging to identify 3 final state particles with S1 only.

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# **Model-independent Reach**

Non-trivial to find appropriate parameterizations for providing model-independent reaches due to many parameters involved in the model

 $\Box$  Number of signal events  $N_{sig}$  is

$$N_{\rm sig} = \sigma \cdot \mathcal{F} \cdot A \cdot t_{\rm exp} \cdot N_e$$

- $\sigma$ : scattering cross section between  $\chi_1$  and (target) electron
- $\mathcal{F}$ : flux of incoming (boosted)  $\chi_1$
- A: acceptance
- *t*<sub>exp</sub>: exposure time
- *N<sub>e</sub>*: total number of target electrons

## Controllable!

# **Model-independent Reach: Displaced Vertex**

□ Acceptance determined by the distance between the primary (ER) and the secondary vertices
 ⇒ (relatively) conservative limit to require two correlated vertices in the fiducial volumes
 (also to be distinguished from elastic scattering)



Evaluated under the assumption of cumulatively isotropic  $\chi_1$  flux

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 ⇒ (relatively) conservative limit to require two correlated vertices in the fiducial volumes
 (also to be distinguished from elastic scattering)

**10<sup>4</sup>** Xe1T-34d 90% C.L. with *Xe*1T–1yr *Xe*1T–5yr zero background **10<sup>3</sup>** DEAP-4d  $\times 10^{-37} [s^{-1}]$ DEAP-1yr  $\sigma \cdot \mathcal{F} \ge \frac{2.5}{A(\ell_{\text{lab}}) \cdot t_{\exp} \cdot N_e}$ DEAP-5vr  $10^{2}$ LZ-1vr LZ = 5vrCalculable given **10**<sup>1</sup> a detector Ţ  $10^{0}$ Evaluated under the assumption  $\mathcal{F}$  for ref1 and ref2 evaluated with of cumulatively isotropic  $\chi_1$  flux  $\langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1}$  being 5 × 10<sup>-26</sup> cm<sup>3</sup>s<sup>-1</sup>  $10^{-1}$ 10 100 1000  $\ell_{lab}$  different event-by-event, so taking  $\ell_{lab}^{max}$  for  $\ell_{\text{lab}}^{\text{max}}$  [cm] more conservative limit

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## Model-independent Reach: "Prompt" Decay

□ No measurable/appreciable displaced vertex  $\Rightarrow$  *A*  $\approx$  1, limit relevant to signals with overlaid vertices or elastic scattering signals

$$\sigma \geq \frac{2.3}{\mathcal{F} \cdot A \cdot t_{\exp} \cdot N_e} \text{ with}$$
$$\mathcal{F} \sim \frac{\langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1}}{m_0^2}$$
set to be 5 × 10<sup>-26</sup> cm<sup>3</sup>s<sup>-1</sup>

### Model-independent Reach: "Prompt" Decay

□ No measurable/appreciable displaced vertex  $\Rightarrow$  *A*  $\approx$  1, limit relevant to signals with overlaid vertices or elastic scattering signals



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### **Dark Photon Parameter Space: Invisible X Decay**

 $\Box$  Case study 1: mass spectra for which dark photon decays into DM pairs, i.e.,  $m_X > 2m_1$ 

□ Same selection criteria imposed

Preliminary  $m_1 = 2$  MeV,  $\gamma_1 = 20$ ,  $m_X > 2m_1$  $m_1 = 2$  MeV,  $\gamma_1 = 20, m_X > 2m$  $\epsilon$  $\epsilon$ 10<sup>-3</sup> 10<sup>-3</sup>  $10^{-4}$  $10^{-4}$ 10<sup>-5</sup> 10<sup>-5</sup> Xenon1T 1y,  $\delta m = 0$  $\delta m = 0$ Xenon1T 1y,  $\delta m = 4$  MeV Xenon1T 1y Xenon1T 5y DEAP3600 1y,  $\delta m = 0$ 10-6 10<sup>-6</sup> DEAP3600 1y DEAP3600 1y,  $\delta m = 2$  MeV DEAP3600 5y DEAP3600 1y, single signal LZ 5y Elastic scattering only Elastic vs. inelastic  $\delta m = 2 \text{ MeV}$ 10-7 10-7 10<sup>-3</sup> 10-2 10<sup>-2</sup>  $10^{-3}$  $10^{-1}$  $10^{-1}$  $m_X$  [GeV]  $m_X$  [GeV]

Caused by the position resolution of 6.5 cm at DEAP

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### Dark Photon Parameter Space: Visible X decay

 $\Box$  Case study 2: mass spectra for which dark photon decays into lepton pairs, i.e.,  $m_X < 2m_1$ 

□ Same selection criteria imposed

Preliminary  $m_1 = 20 \text{ MeV}, \gamma_1 = 100, m_X < 2m_1$  $m_1 = 20 \text{ MeV}, \gamma_1 = 100, m_X < 2m_1$  $\epsilon$  $\epsilon$  $10^{-3}$ 10<sup>-3</sup> 10<sup>-4</sup>  $10^{-4}$ Xenon1T 1y,  $\delta m = 0$  $\delta m = 0$ Xenon1T 1y,  $\delta m = 20$  MeV Xenon1T 1y Xenon1T 5y DEAP3600 1y,  $\delta m = 0$ DEAP3600 1y DEAP3600 1y,  $\delta m = 20$  MeV DEAP3600 5y DEAP3600 1y, single signal 10<sup>-5</sup> LZ 5y **Elastic vs. inelastic**  $\delta m = 20 \text{ MeV}$ attering only 10<sup>-5</sup> 10<sup>-2</sup>  $10^{-3}$  $10^{-2}$  $10^{-3}$  $10^{-1}$  $10^{-1}$  $m_X$  [GeV]  $m_X$  [GeV]

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## Conclusions

- Boosted light dark matter searches are **promising**.
- Conventional dark matter direct detection experiments possess sensitivities to MeV-range (heaviest light?) DM.
- They can provide an alternative avenue to probe dark photon parameter space.







### **Boosted DM from the Sky: Semi-annihilation**

In DM models where relevant DM is stabilized by e.g., Z<sub>3</sub> symmetry, one may have a process like



□ Under the circumstance in which the mass of SM here is lighter (i.e.,  $m_A > m_{SM}$ ), the outgoing  $\chi_A$  can be boosted and its boost factor is given by

$$\gamma_A = \frac{5m_A^2 - m_{\rm SM}^2}{4m_A^2}$$

### **Boosted DM Signal Detection**

#### [LUX Collaboration (2017)]



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# **Backgrounds for Xenon1T**

**Table 2** Summary of the sources contributing to the background of XENON1T in a fiducial target of 1.0 t and a NR energy region from 4 to 50 keV (corresponding to 1 to 12 keV ER equivalent). The expected rates are taken from the Monte Carlo simulation-based study [18] and assume no ER rejection. CNNS stands for "coherent neutrino nucleus scattering".

Background Source	Туре	Rate $[(t \times y)^{-1}]$	Mitigation Approach
$222 \operatorname{Rn} \left(10 \mu \mathrm{Bq/kg}\right)$	ER	620	material selected for low Rn-emanation; ER rejection
solar pp- and <sup>7</sup> Be-neutrinos	ER	36	ER rejection
$^{85}$ Kr (0.2 ppt of <sup>nat</sup> Kr)	ER	31	cryogenic distillation; ER rejection
$2\nu\beta\beta$ of <sup>136</sup> Xe	ER	9	ER rejection
Material radioactivity	ER	30	material selection; ER and multiple scatter rejection; fiducialization
Radiogenic neutrons	NR	0.55	material selection; multiple scatter rejection; fiducialization
CNNS (mainly solar <sup>8</sup> B-neutrinos)	NR	0.6	_
Muon-induced neutrons	NR	< 0.01	active Cherenkov veto [43]; multiple scatter rejection; fiducialization

[Xenon Collaboration (2017)]

### All are smaller than ~100 keV, hence irrelevant to our signals