

HW 6 Sol

$$5.1) a) e^{-i x_0 p / \hbar} \psi = e^{-x_0 \partial_x} \psi$$

$$= \sum_n \frac{(-x_0)^n}{n!} \psi^{(n)}(x)$$

\swarrow n th derivative

compare Taylor Series $\psi(x-x_0) = \sum_n \frac{(-x_0)^n}{n!} \psi^{(n)}(x)$

$\rightarrow = \psi(x-x_0)$

Alternate

$$e^{-i x_0 p / \hbar} |\psi\rangle = \int \frac{dp}{(2\pi\hbar)^{1/2}} e^{-i x_0 p / \hbar} e^{i p x / \hbar} \varphi(p)$$

$$= \int \frac{dp}{(2\pi\hbar)^{1/2}} e^{i p (x-x_0) / \hbar} \varphi(p)$$

$$= \psi(x-x_0)$$

$$b) e^{-i \phi L_z / \hbar} \psi(r, \theta) = \int_{-\pi}^{\pi} e^{-i \phi \partial_\theta} \psi(r, \theta)$$

$$= \sum_n \frac{(-\phi)^n}{n!} \partial_\theta^n \psi(r, \theta) = \psi(r, \theta - \phi)$$

or, $\psi(r, \theta) = \sum_m \psi_m(r) e^{i m \theta}$

$$e^{-i \phi \partial_\theta} e^{i m \theta} = \sum_n \frac{(-\phi)^n}{n!} (i m)^n e^{i m \theta} = e^{-i m \phi} e^{i m \theta}$$

$$= e^{i m (\theta - \phi)}$$

$$\therefore e^{-i \phi L_z / \hbar} \psi(r, \theta) = \sum_m \psi_m(r) e^{i m (\theta - \phi)} = \psi(r, \theta - \phi)$$

$$5.2) \quad e^{-iH(t-t_0)/\hbar} \psi(t_0) = \sum_n \frac{(-iH(t-t_0)/\hbar)^n}{n!} \psi(t_0)$$

$$H\psi = i\hbar \frac{d}{dt} \psi$$

$$\Rightarrow = \sum_n \frac{(t-t_0)^n}{n!} \frac{d^n \psi(t_0)}{dt^n} = \psi(t)$$

$$\text{or, } \psi(t) = \sum_n e^{-iE_n(t-t_0)/\hbar} |n\rangle \langle n|\psi(t_0)\rangle$$

$$\left(\text{where } H|n\rangle = E_n|n\rangle \right)$$

$$= \sum_n e^{-iH(t-t_0)/\hbar} |n\rangle \langle n|\psi(t_0)\rangle$$

$$= e^{-iH(t-t_0)/\hbar} \psi(t_0)$$

$$\text{or, } \psi(t_0) = \int \frac{d\omega}{2\pi} \tilde{\psi}(\omega) e^{-i\omega t_0}$$

$$e^{-iH(t-t_0)/\hbar} e^{-i\omega t_0} = e^{(t-t_0)d/dt} e^{-i\omega t_0}$$

$$= \sum_n \frac{(t-t_0)^n}{n!} (-i\omega)^n e^{-i\omega t_0}$$

$$= e^{-i\omega t}$$

$$\text{so } e^{-iH(t-t_0)/\hbar} \psi(t_0) = \int \frac{d\omega}{2\pi} \tilde{\psi}(\omega) e^{-i\omega t}$$

$$= \psi(t)$$

$$b) \quad U^\dagger \cdot U = e^{iH(t-t_0)/\hbar} e^{-iH(t-t_0)/\hbar} = 1$$

$$\text{so } U^\dagger = U^{-1}$$

5.3)

$$a(t) = \langle \psi(t) | A | \psi(t) \rangle$$

$$= \langle \psi(0) | U^\dagger(t) A U(t) | \psi(0) \rangle$$

$$= \langle \psi(0) | A(t) | \psi(0) \rangle$$

$$i\hbar \frac{dA}{dt} = i\hbar \underbrace{\left(\frac{d}{dt} U^\dagger\right)}_{\frac{iH U^\dagger}{\hbar}} A U + i\hbar U^\dagger A \underbrace{\frac{dU}{dt}}_{-\frac{iU H}{\hbar}}$$

$$= -H A(t) + A(t) H$$

$$= [A, H]$$

5.4) $\langle \psi_2 | H | \psi_1 \rangle = E_2 \langle \psi_2 | \psi_1 \rangle = E_1 \langle \psi_2 | \psi_1 \rangle$

$$\therefore (E_2 - E_1) \langle \psi_2 | \psi_1 \rangle = 0 \Rightarrow \langle \psi_2 | \psi_1 \rangle = 0$$

b) $\langle \frac{\psi_1 - \psi_2}{\sqrt{2}} | H | \frac{\psi_1 - \psi_2}{\sqrt{2}} \rangle = \frac{E_1}{2} + \frac{E_2}{2} = \frac{E_1 + E_2}{2}$

$$\langle H^2 \rangle = \frac{E_1^2}{2} + \frac{E_2^2}{2}$$

$$\Delta E = \left(\frac{E_1^2}{2} + \frac{E_2^2}{2} - \left(\frac{E_1 + E_2}{2} \right)^2 \right)^{1/2} = \left(\frac{(E_1 - E_2)^2}{4} \right)^{1/2} = \frac{E_1 - E_2}{2}$$

c) $|\psi(t)\rangle = \frac{1}{\sqrt{2}} \left(e^{-iE_1 t/\hbar} |\psi_1\rangle - e^{-iE_2 t/\hbar} |\psi_2\rangle \right)$

d) $A = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \det \begin{pmatrix} -\lambda & 1 \\ 1 & -\lambda \end{pmatrix} = \lambda^2 - 1 \Rightarrow \lambda = \pm 1$

e) Eigenvectors $|+1\rangle$ $(-1)\langle\psi_1|+1\rangle + \langle\psi_2|+1\rangle = 0$

o. $|+1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$

$| -1 \rangle$ $\langle\psi_1|-1\rangle + \langle\psi_2|-1\rangle = 0$

$| -1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$

f) $P_{-1} = |\langle -1|\psi\rangle|^2$

$$\begin{aligned} \langle -1|\psi\rangle &= \frac{1}{\sqrt{2}} e^{-iE_1 t/\hbar} \underbrace{\langle -1|\psi_1\rangle}_{1/\sqrt{2}} - \frac{1}{\sqrt{2}} e^{-iE_2 t/\hbar} \underbrace{\langle -1|\psi_2\rangle}_{-1/\sqrt{2}} \\ &= e^{-i\frac{(E_1+E_2)}{2\hbar}t} \left(\frac{1}{2} e^{-i\omega t} + \frac{1}{2} e^{i\omega t} \right) \\ &= e^{-i\dots} \cos \omega t \end{aligned}$$

o. $P_{-1} = \cos^2 \omega t$

5) $H = \text{diag}(0, \Delta, 0) + V$

$= \begin{pmatrix} 0 & \epsilon & 0 \\ \epsilon^* & \Delta & \epsilon \\ 0 & \epsilon^* & 0 \end{pmatrix}$ ~~or~~

Mathematica let $\tan 2\theta = \frac{\sqrt{8}|\epsilon|}{\Delta}$, $\Omega = \sqrt{8|\epsilon|^2 + \Delta^2}$

$E=0$, $|0\rangle = \begin{pmatrix} -1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix}$ $E = -\Omega \sin^2 \theta$, $|-\rangle = \begin{pmatrix} \cos \theta / \sqrt{2} \\ -e^{-i\phi} \sin \theta \\ \cos \theta / \sqrt{2} \end{pmatrix}$

$E = \Omega \cos^2 \theta$, $|+\rangle = \begin{pmatrix} \sin \theta / \sqrt{2} \\ e^{-i\phi} \cos \theta \\ \sin \theta / \sqrt{2} \end{pmatrix}$

dark state is $|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$

$$V|0\rangle = \begin{pmatrix} 0 & \epsilon & 0 \\ \epsilon^* & 0 & \epsilon \\ 0 & \epsilon^* & 0 \end{pmatrix} \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} = 0$$

in that state, there is no interaction with the light!

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In[235]:= {eval, kets} = Eigensystem[ $\begin{pmatrix} 0 & e e^{i\phi} & 0 \\ e e^{-i\phi} & \Delta & e e^{-i\phi} \\ 0 & e e^{i\phi} & 0 \end{pmatrix}$  /. {e ->  $\frac{\sin[2\theta]}{\sqrt{8}}$   $\Omega$ ,  $\Delta$  ->  $\cos[2\theta] \Omega$ }] //
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TrigFactor // Simplify
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Out[235]= {{0, - $\Omega \sin[\theta]^2$ ,  $\Omega \cos[\theta]^2$ },  
{{-1, 0, 1}, {1, - $\sqrt{2} e^{-i\phi} \tan[\theta]$ , 1}, {1,  $\sqrt{2} e^{-i\phi} \cot[\theta]$ , 1}}}
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note: *Mathematica* gives the eigenvectors as row vectors. So the individual kets are kets[[1]], kets[[2]] etc. The bras are conj[kets[[1]]] and so on. When *Mathematica* calculates the eigenvectors and eigenvalues of a numerical matrix, it gives properly normalized eigenvectors. However, for an analytical matrix they are not normalized so you have to take care of the normalization yourself.

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In[236]:= norm = Table[conj[kets[[i]]].kets[[i]], {i, 3}] //  
Simplify (*Calculate the norms of the kets*)
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Out[236]= {2, 2 Sec[ $\theta$ ]2, 2 Csc[ $\theta$ ]2}
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In[237]:= kets =  $\frac{\text{kets}}{\sqrt{\text{norm}}}$  // PowerExpand // Simplify
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Out[237]= {{- $\frac{1}{\sqrt{2}}$ , 0,  $\frac{1}{\sqrt{2}}$ }, { $\frac{\cos[\theta]}{\sqrt{2}}$ , - $e^{-i\phi} \sin[\theta]$ ,  $\frac{\cos[\theta]}{\sqrt{2}}$ }, { $\frac{\sin[\theta]}{\sqrt{2}}$ ,  $e^{-i\phi} \cos[\theta]$ ,  $\frac{\sin[\theta]}{\sqrt{2}}$ }}
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